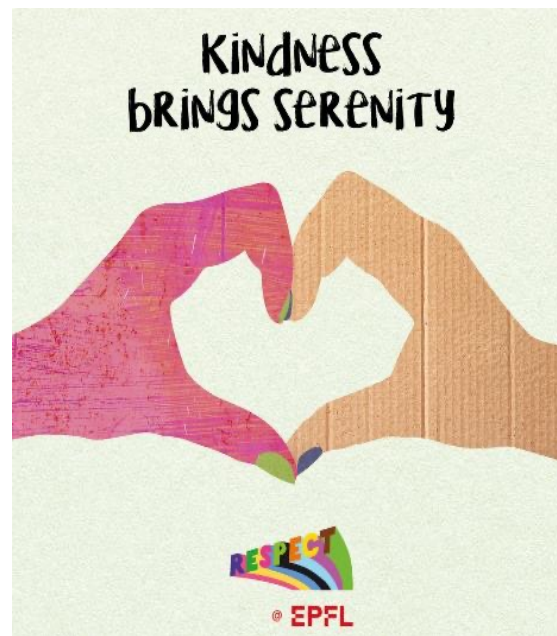
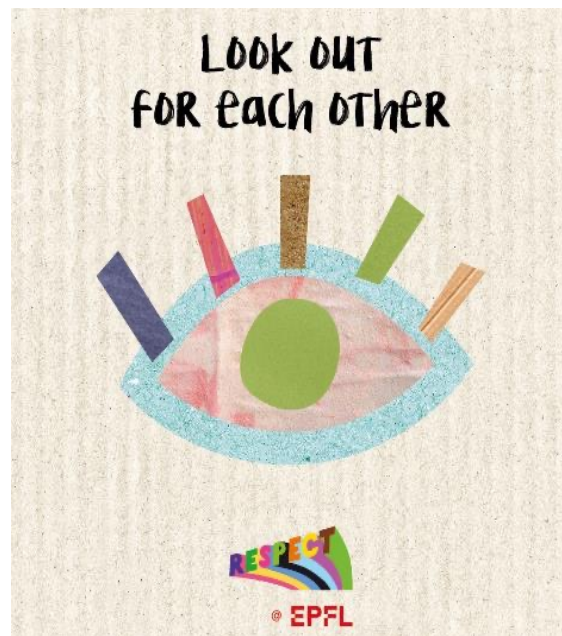
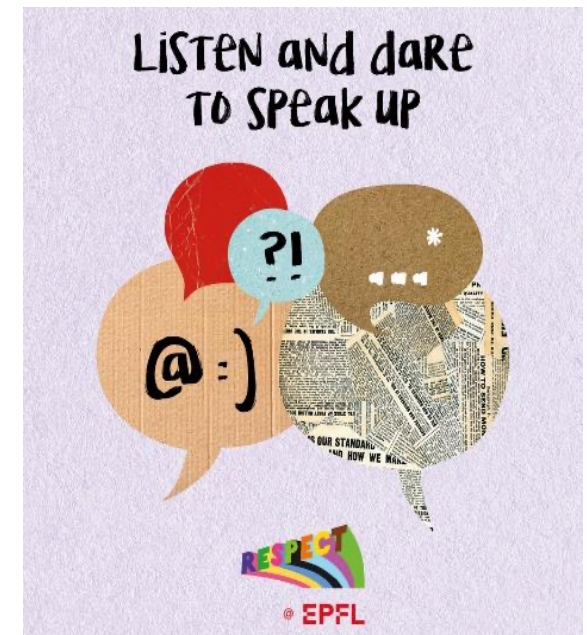
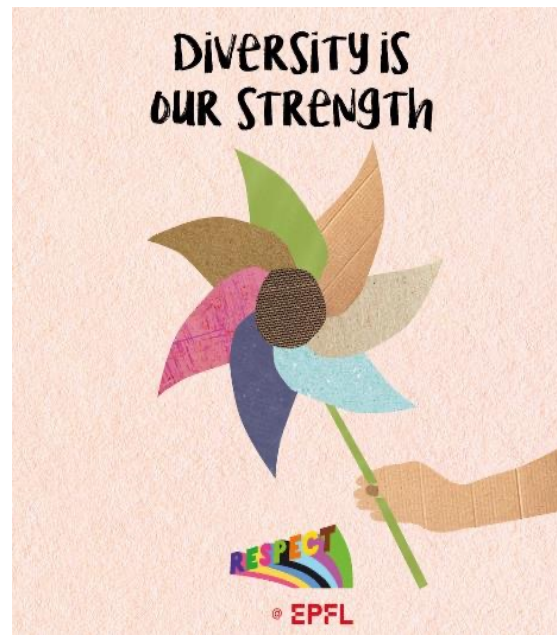
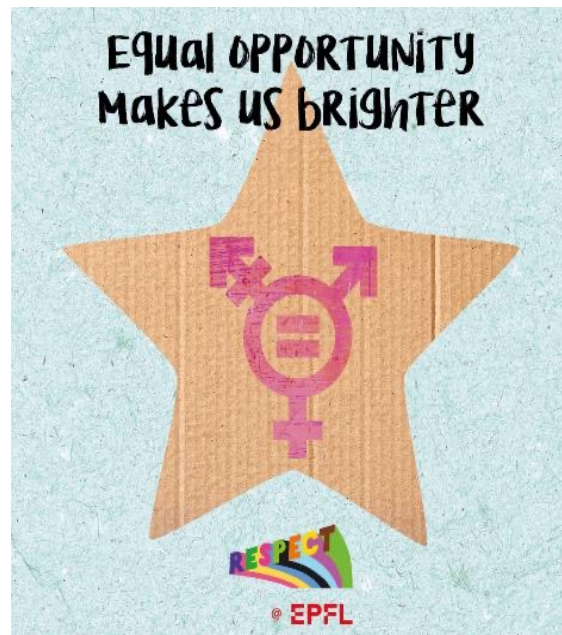




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Génie Electrique et Electronique
Master program
Prof. Elisa Matioli

EE-557 Semiconductor devices I

Doping and carrier concentration

Outline of the lecture

1. Carrier concentration
2. Determining Fermi level

Read Chapter 2 of the reference book (on moodle)

References:

- J. A. del Alamo, course materials for 6.720J Integrated Microelectronic Devices, Spring 2007. MIT OpenCourseWare (<http://ocw.mit.edu/>)

Crystal structure: Definitions

Unit cell: Defines the symmetry and structure of the entire lattice

Bravais lattices: describe the geometric arrangement of the lattice points

a is the lattice constant.

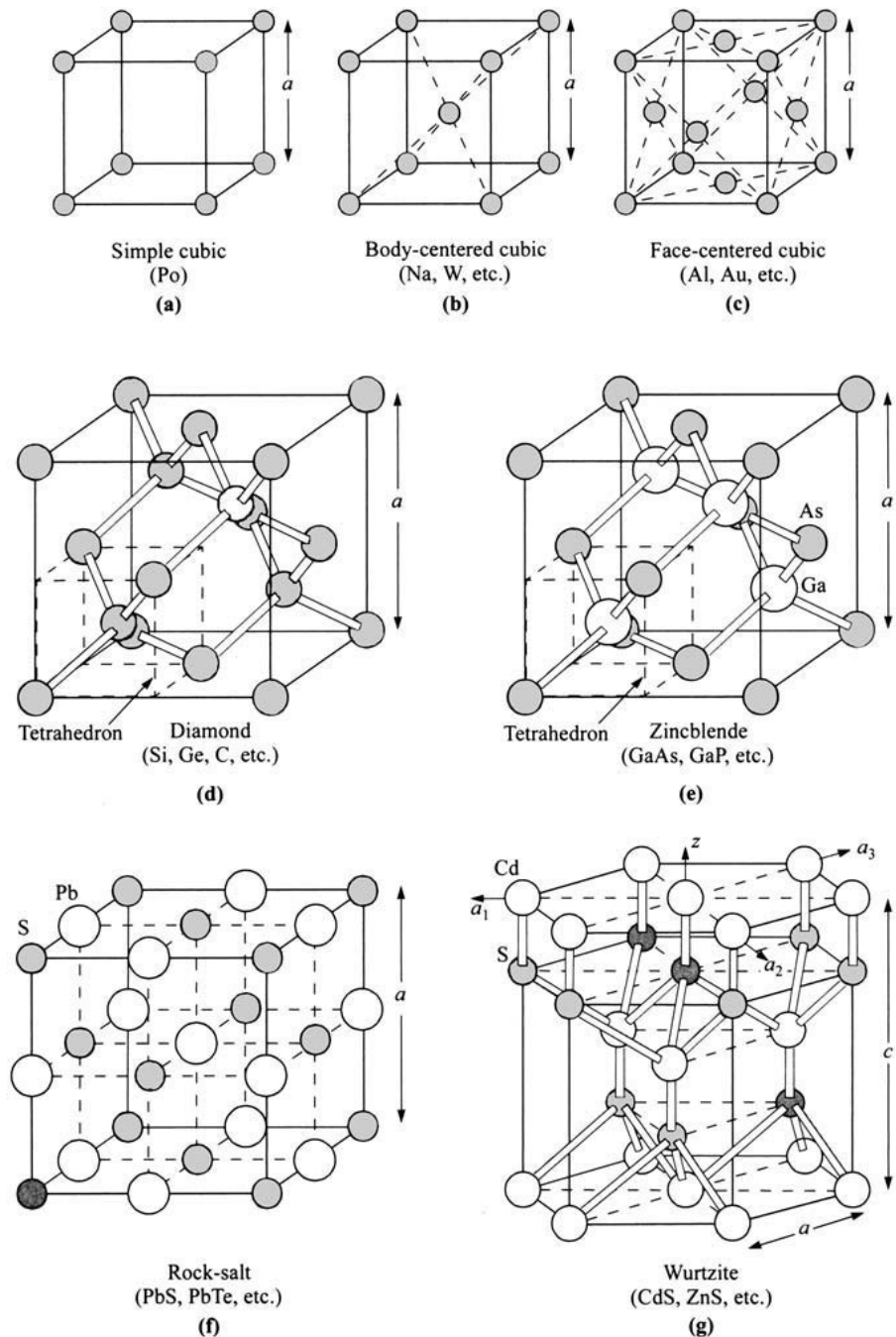
Crystal structure and symmetry play a critical role in determining many physical properties, such as:

- **Cleavage:**
Check out this video:
<https://www.youtube.com/watch?v=IRoIXjlcBQ>
- **Mechanical properties** (elastic compliance, stiffness)
- **Electronic band structure:** Band gap
- **Optical transparency**
- **Thermal properties**
- **Polarization fields**

GaN: most common in wurtzite structure

Zincblende structure: III-V compound semiconductors:
GaAs, GaP, etc:

Important for optoelectronics and high-speed ICs



How Energy Bands and Energy Gap are calculated?

Energy-momentum relationship: characterizes the band structure

- Important for the interactions with photons and phonons

Schrödinger's equation:
$$\left[-\frac{\hbar^2}{2m^*} \nabla^2 + V(\mathbf{r}) \right] \psi(\mathbf{r}, \mathbf{k}) = E(\mathbf{k}) \psi(\mathbf{r}, \mathbf{k})$$

Bloch function
$$\psi(\mathbf{r}, \mathbf{k}) = \exp(j\mathbf{k} \cdot \mathbf{r}) U_b(\mathbf{r}, \mathbf{k})$$

$\psi(\mathbf{r}, \mathbf{k})$ and $U_b(\mathbf{r}, \mathbf{k})$ are periodic in $\mathbf{r}$ in real space

Thus:

$$\begin{aligned} \psi(\mathbf{r} + \mathbf{R}, \mathbf{k}) &= \exp[j\mathbf{k} \cdot (\mathbf{r} + \mathbf{R})] U_b(\mathbf{r} + \mathbf{R}, \mathbf{k}) \\ &= \exp(j\mathbf{k} \cdot \mathbf{r}) \exp(j\mathbf{k} \cdot \mathbf{R}) U_b(\mathbf{r}, \mathbf{k}) \end{aligned}$$

$\mathbf{k} \cdot \mathbf{R}$ is a multiple of 2π .

How Energy Bands and Energy Gap are calculated?

In 1D – simple case

Consequences:

- In 1D: only $k = 2\pi/L$ are allowed (where a is the real space period)
- $E(k)$ is periodic in k -space: $E(k) = E(k+G)$
- It is sufficient to define k in a primitive cell, which is defined by the Brillouin zone: $-\pi/a, \pi/a$
- Entire band structures need only to be calculated within the Brillouin zone.

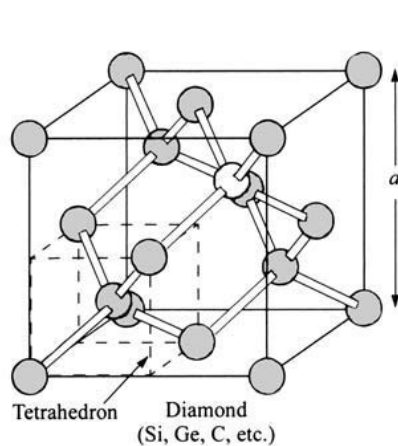
In 3D - Reciprocal space

Defines the principal symmetry points :

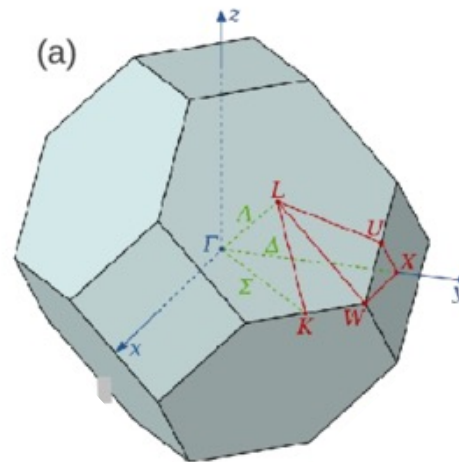
Gamma = (0; 0; 0) is the origin

X = (1; 0; 0) + 6 equivalent points,

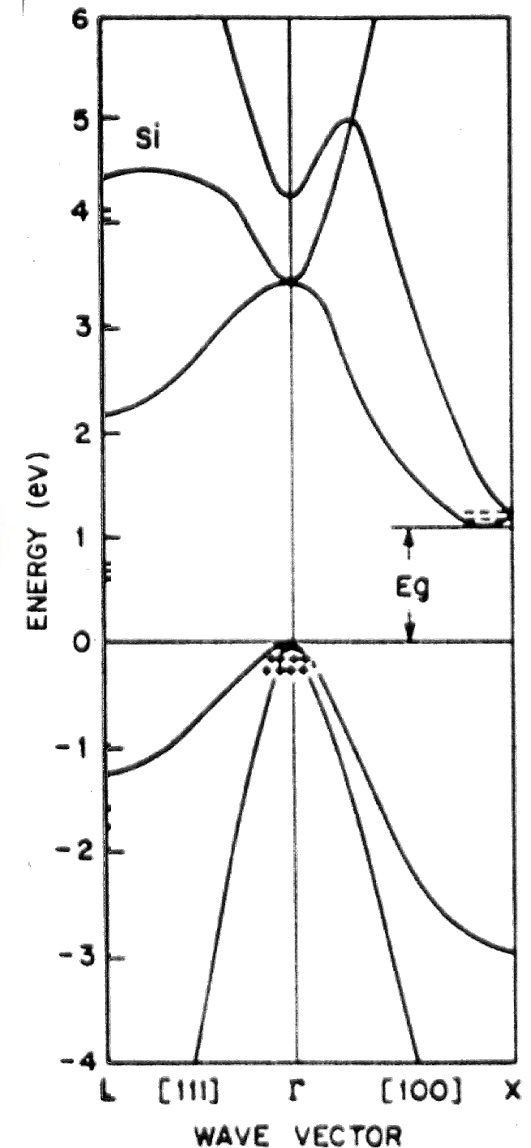
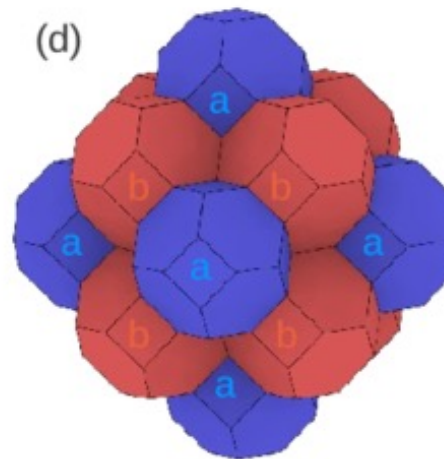
L = (1; 1; 1) + 8 equivalent points



Real space



Reciprocal space



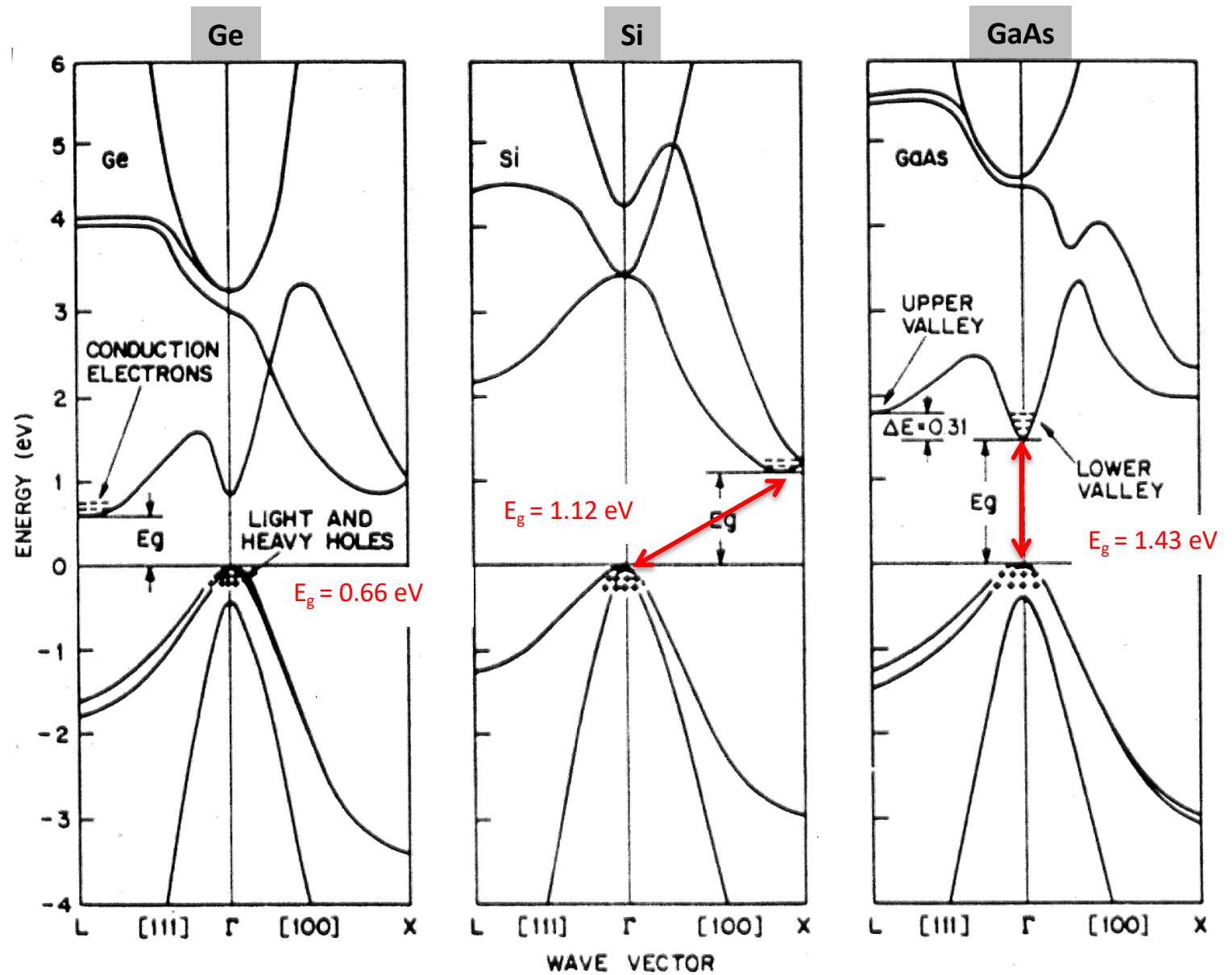
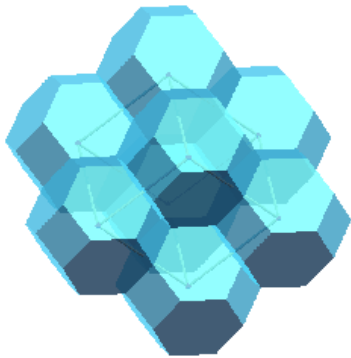
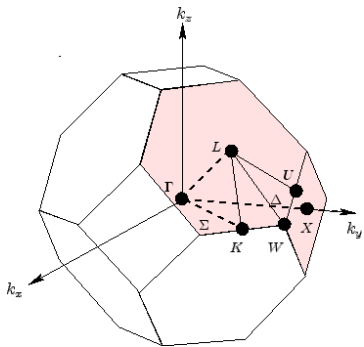
Fundamental to define the band structure of a semiconductor

The reciprocal lattice plays a fundamental role in most analytic studies (properties related to momentum) of periodic structures:

- Energy-momentum relationship
- Theory of diffraction: X-Ray diffraction

Energy band-gap

Reciprocal space directions:



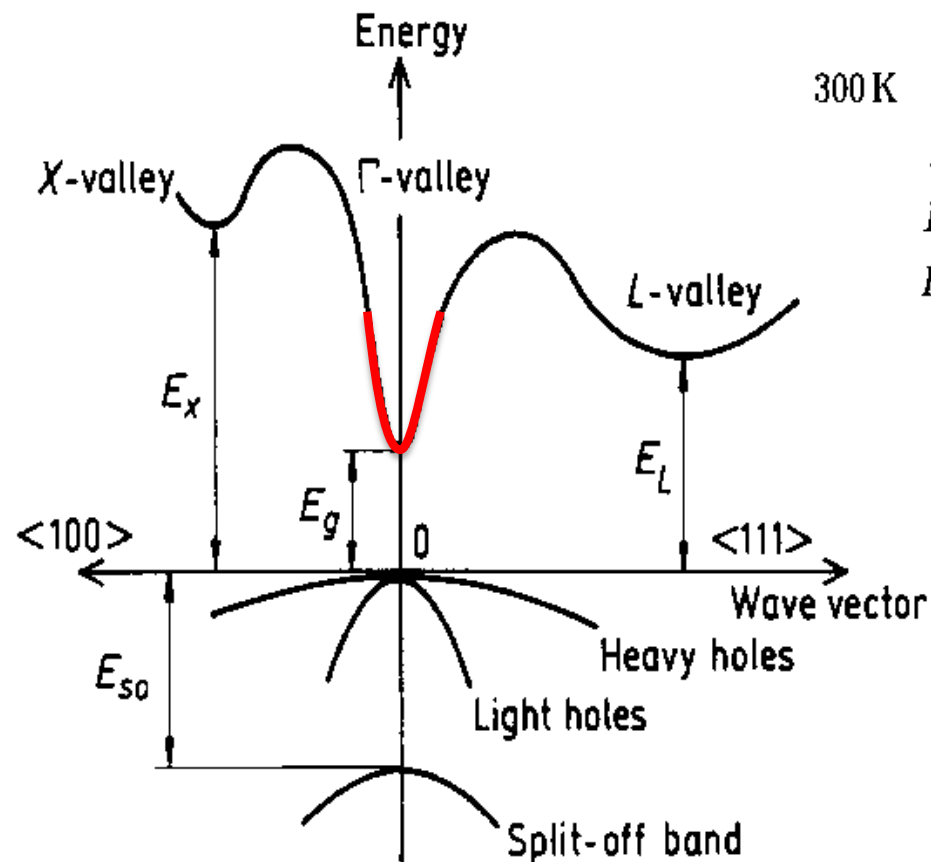
Indirect band gap

direct band gap

What do we learn from band diagrams

Band diagram: We focus on minimum of the bands
(quadratic region)

GaAs



What do we learn from it:

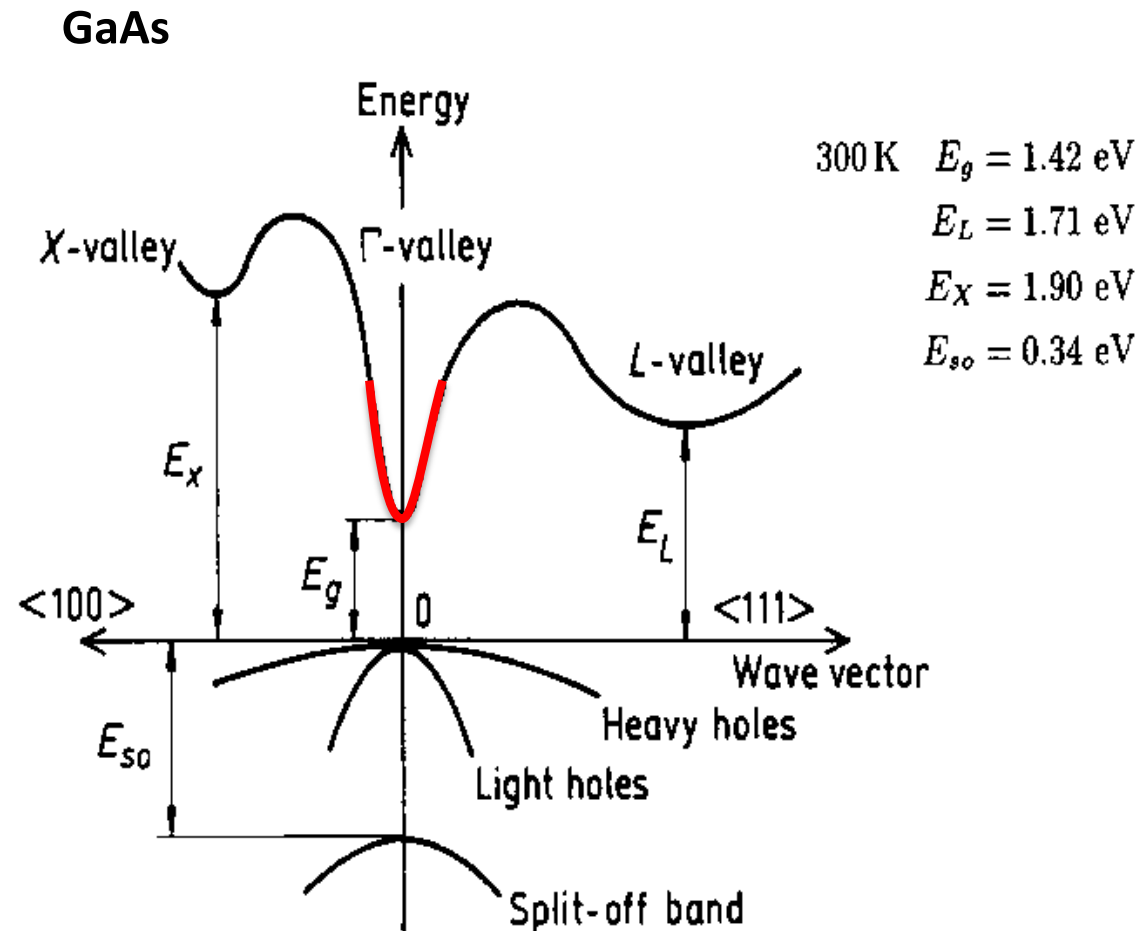
Energy-momentum relationship near band edges

1. Allowed and forbidden states
2. Group velocity: slope of the bands
3. Electron/hole mass: curvature of bands

What do we learn from band diagrams

Effective mass

Band diagram: We focus on minimum of the bands (quadratic region)



Energy-momentum relationship near band edges

$$E(k) = \frac{p^2}{2m^*} = \frac{\hbar^2 k^2}{2m^*}$$

Effective mass in 1D

$$m_j^* = \hbar^2 \left(\frac{\partial^2 E_j}{\partial k_y^2} \right)^{-1}$$

$\hbar = h/2\pi$: planck's constant = 6.58 eV.s

For example:

Si: $m^* = 1.09 m_0$

GaN: $m^* = 0.2 m_0$

GaAs: $m^* = 0.06 m_0$

Group velocity:

$$v_n(k) = \frac{1}{\hbar} \nabla_k E_n(k)$$

$$v_n(k) = \frac{\hbar k}{m^*} = \frac{p}{m^*}$$

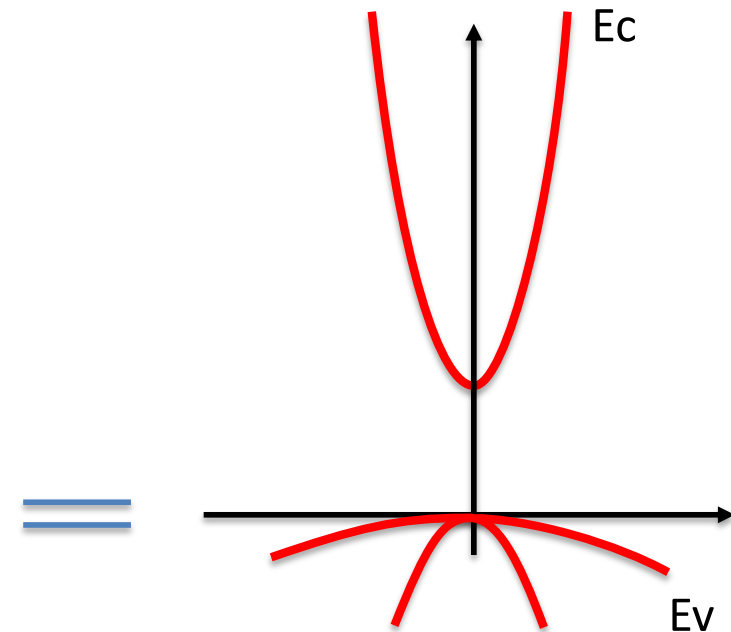
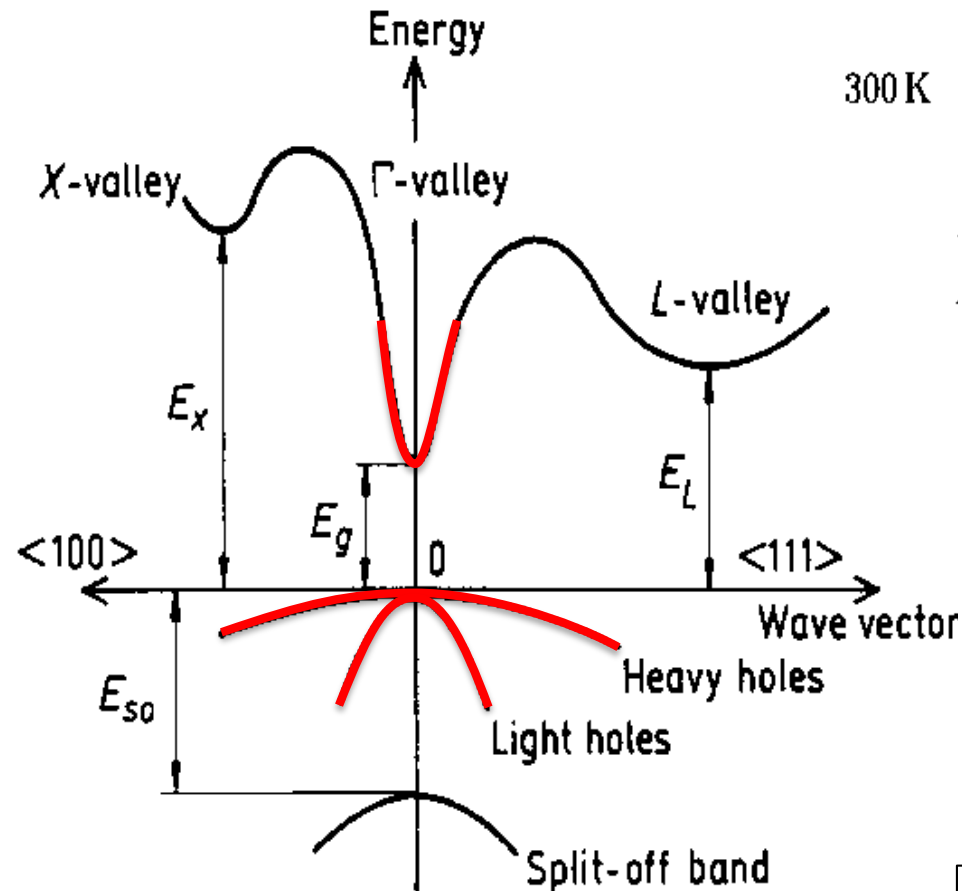
Momentum $p = \hbar k$

- How many carriers are available for conduction?
- How can we increase and control the conductivity in semiconductors?

A few important concepts

How many carriers are available for conduction?

We are interested in analyzing the minimum of the conduction and valence bands



Available electrons in the conduction band:

$$n_o = \int_{E_c}^{\infty} g_c(E) f(E) dE$$

density of electrons

Probability of occupation

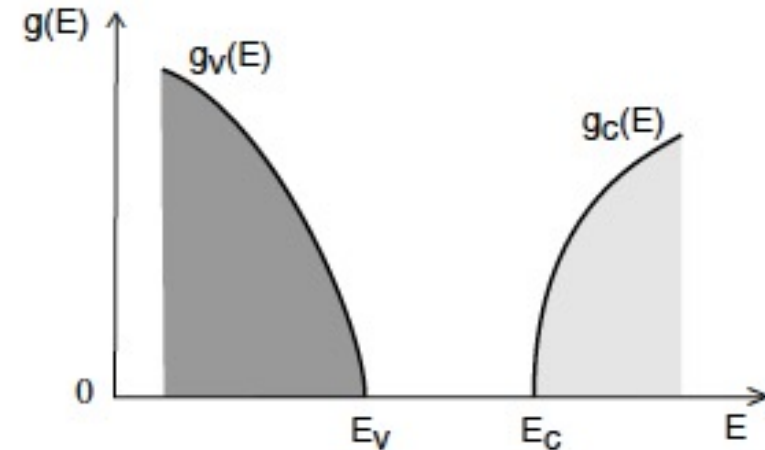
What is the density of electrons $g_c(E)$?

What law regulates the electron occupation of states $f(E)$ as a function of energy and temperature?

$$n_o = \int_{E_c}^{\infty} g_c(E) f(E) dE$$

A few important concepts

Density of states



$$g_c(E) = 4\pi \left(\frac{2m_{de}^*}{h^2} \right)^{3/2} \sqrt{E - E_c} \quad E \geq E_c$$

$$g_v(E) = 4\pi \left(\frac{2m_{dh}^*}{h^2} \right)^{3/2} \sqrt{E_v - E} \quad E \leq E_v$$

m_{de}^* $\equiv$ density of states electron effective mass

m_{dh}^* $\equiv$ density of states hole effective mass

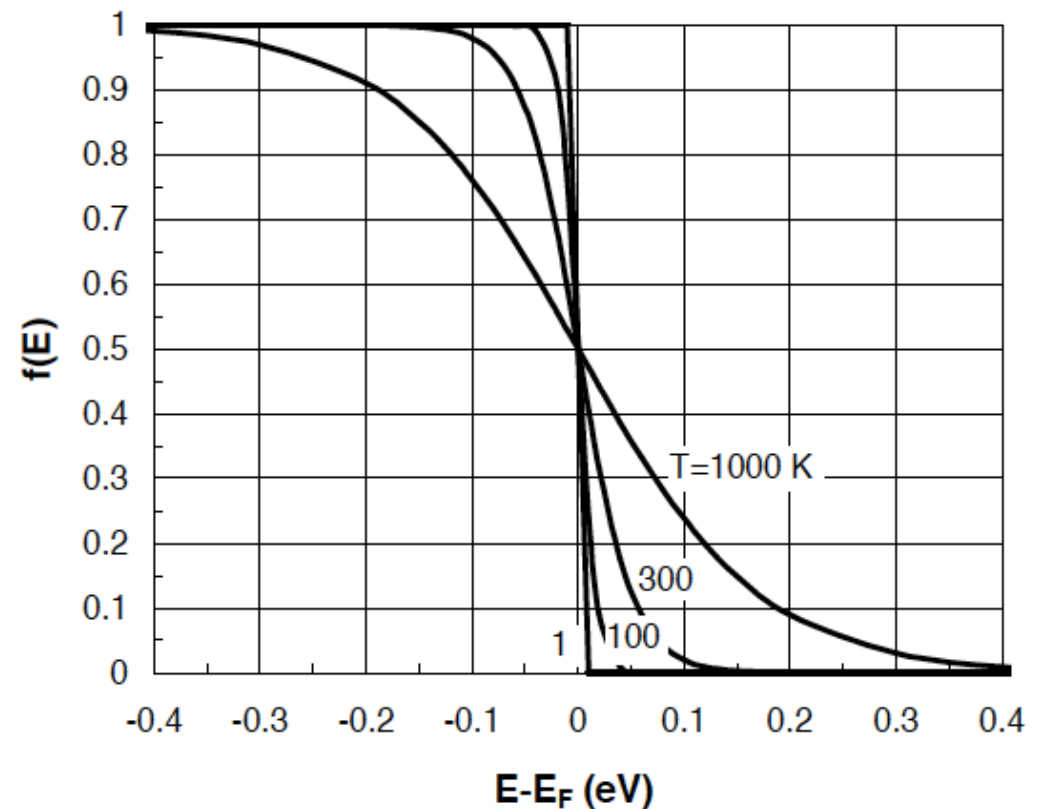
A few important concepts

Density of states

Electron statistics

At finite temperature, state occupation probability by electron determined by **Fermi-Dirac distribution function**:

$$f(E) = \frac{1}{1 + \exp \frac{E - E_F}{kT}}$$



E_F : Fermi energy -energy for which occupation probability is 50%

k : Boltzmann constant = 8.62×10^{-5} eV/K

kT -thermal energy = 25.9 meV at 300 K

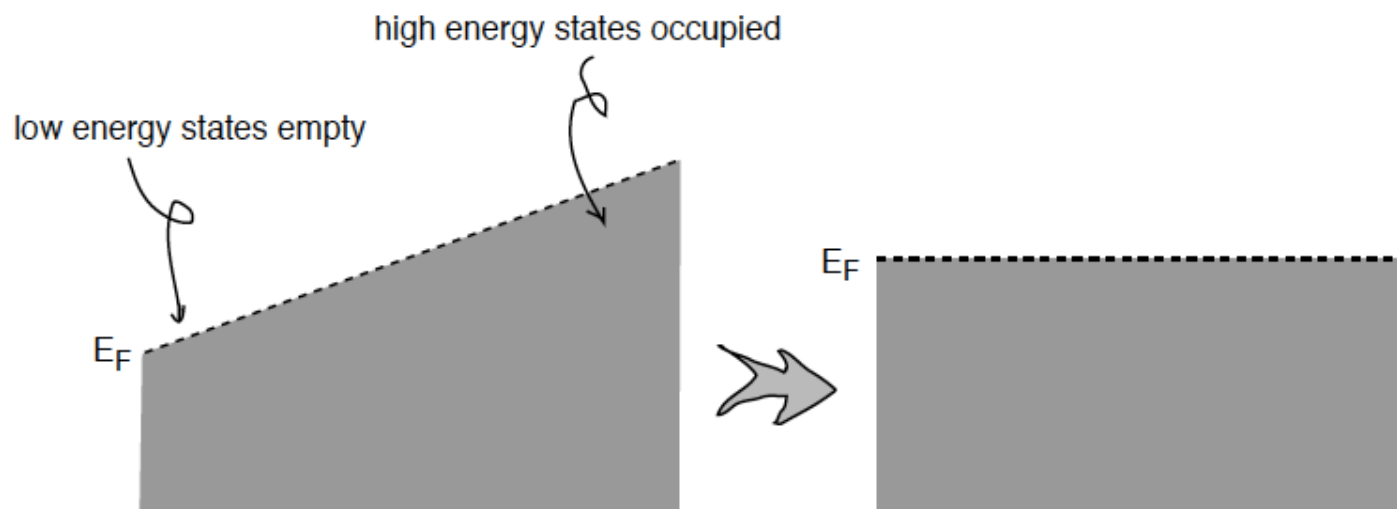
Notion of thermal equilibrium

A particle system is in thermal equilibrium (TE) if:

- it is closed: no energy flows through boundaries of system
- it is in steady-state: time derivatives of all ensemble averages (global and local) are zero

Thermal equilibrium important because all systems evolve towards TE after having been perturbed.

In thermal equilibrium, E_F constant throughout system



How to calculate the carrier concentration? (which is related to the current)

Density of states

$$g_c(E) = 4\pi \left(\frac{2m_{de}^*}{h^2} \right)^{3/2} \sqrt{E - E_c} \quad E \geq E_c$$

$$m_{de} = (m_1^* m_2^* m_3^*)^{1/3}$$

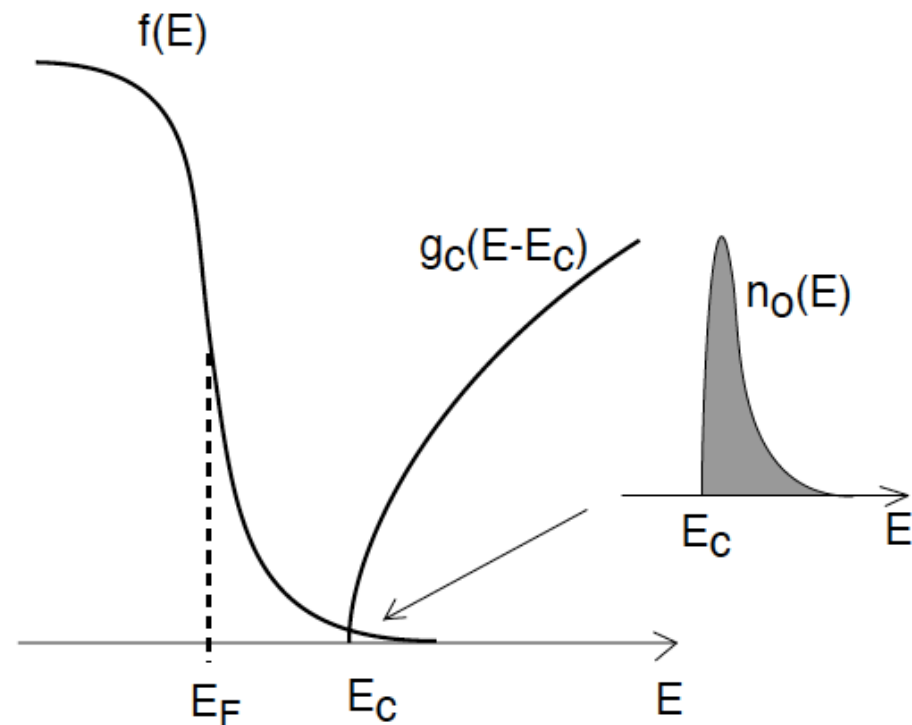
Probability of a state to be occupied

$$f(E) = \frac{1}{1 + \exp \frac{E - E_F}{kT}}$$

X

Integrated for all energies

$$n_o = \int_{E_c}^{\infty} g_c(E) f(E) dE$$

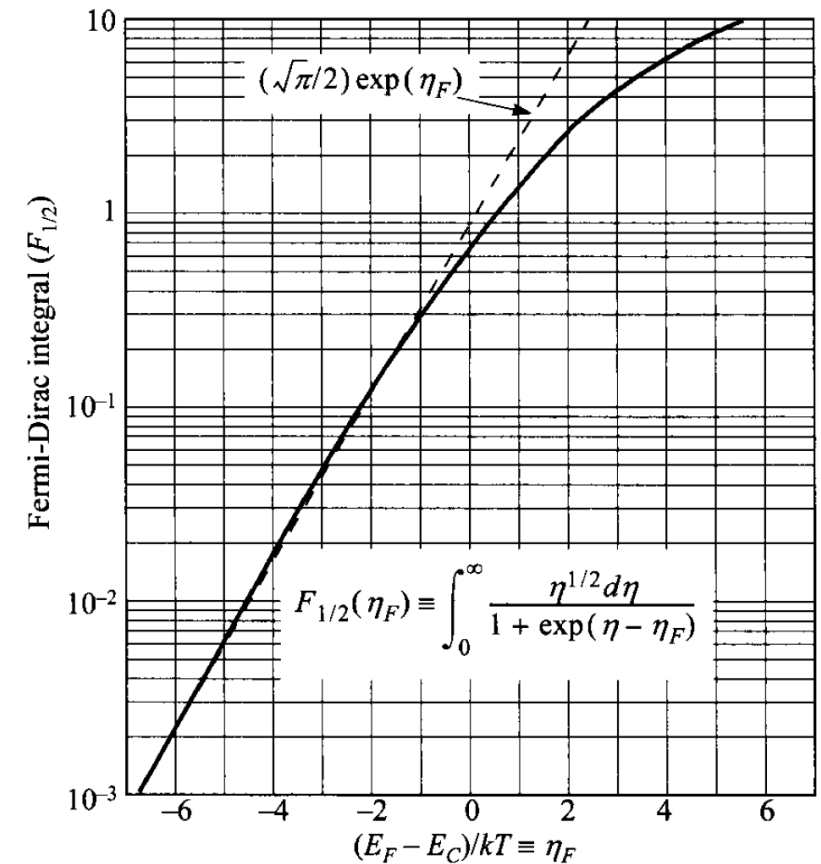


1. What is a simplified expression for the number of electrons n and holes p ?
2. What is the position of the Fermi level?
3. What is the np product?

$$n = N_C \frac{2}{\sqrt{\pi}} F_{1/2} \left(\frac{E_F - E_C}{kT} \right)$$

$$\text{where } N_C \equiv 2 \left(\frac{2 \pi m_{de} kT}{h^2} \right)^{3/2}$$

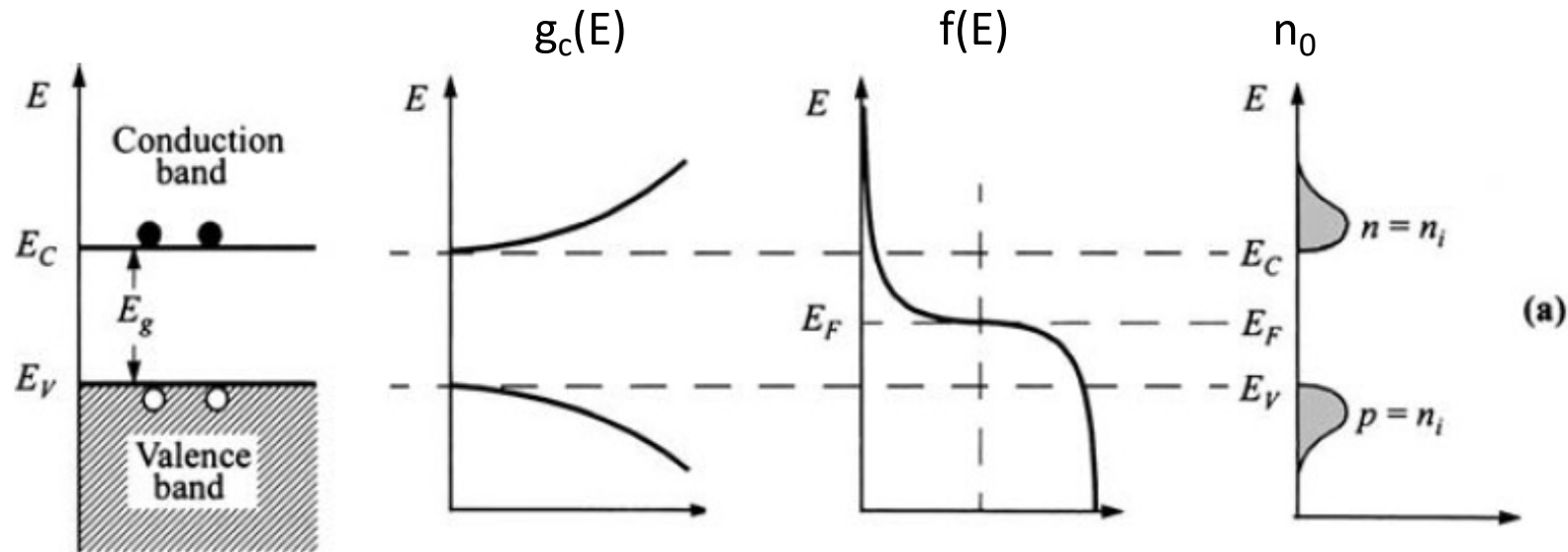
$$\begin{aligned} F_{1/2} \left(\frac{E_F - E_C}{kT} \right) &\equiv F_{1/2}(\eta_F) = \int_{E_C}^{\infty} \frac{[(E - E_C)/kT]^{1/2}}{1 + \exp[(E - E_F)/kT]} \frac{dE}{kT} \\ &= \int_0^{\infty} \frac{\eta^{1/2}}{1 + \exp(\eta - \eta_F)} d\eta \end{aligned}$$



Non-degenerate Semiconductors. By definition, in non-degenerate semiconductors, the doping concentrations are smaller than N_C and the **Fermi levels are more than several kT below E_C** , the Fermi-Dirac integral approaches:

$$F_{1/2} \left(\frac{E_F - E_C}{kT} \right) = \frac{\sqrt{\pi}}{2} \exp \left(- \frac{E_C - E_F}{kT} \right)$$

Intrinsic semiconductor (undoped)



$n_o \equiv$ equilibrium (free) electron concentration in conduction band [cm^{-3}]

$p_o \equiv$ equilibrium hole concentration in valence band [cm^{-3}]

In intrinsic semiconductors:

$n_o = p_o = n_i$ $n_i \equiv$ intrinsic carrier concentration [cm^{-3}]

Carrier concentration

Intrinsic semiconductor (undoped)

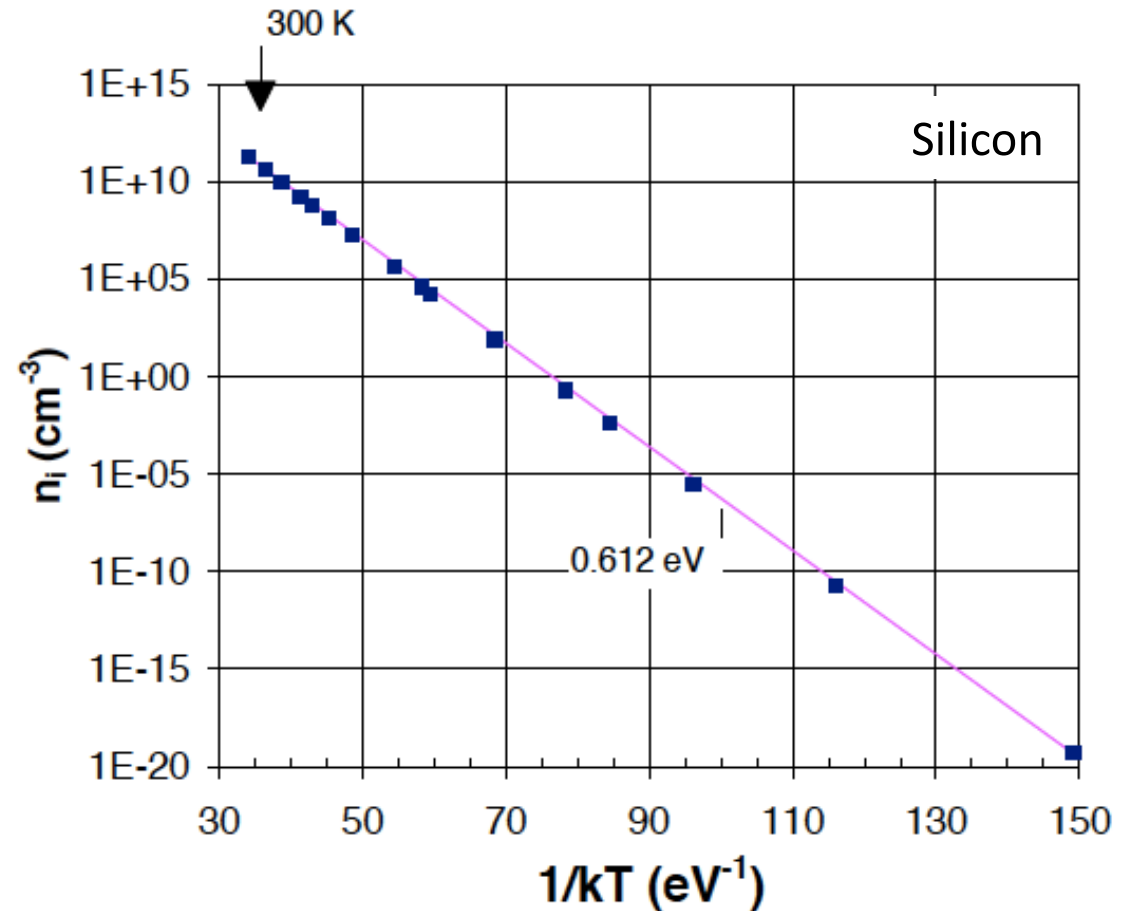
First important result:

$$n_o = p_o = n_i$$

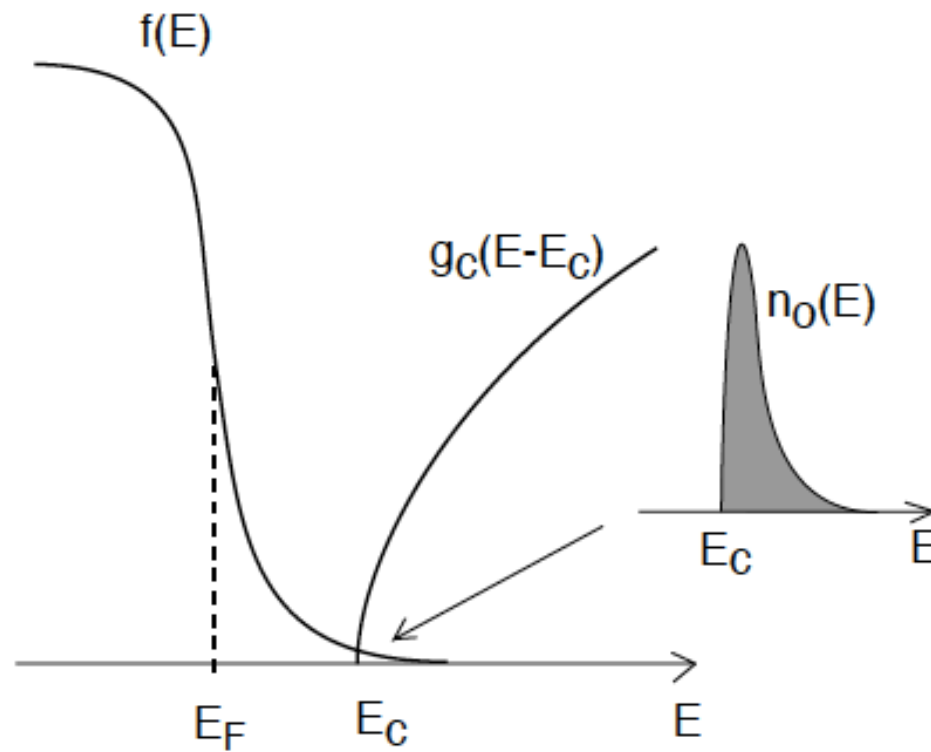
$n_i \equiv$ intrinsic carrier concentration [cm^{-3}]

A second important result:

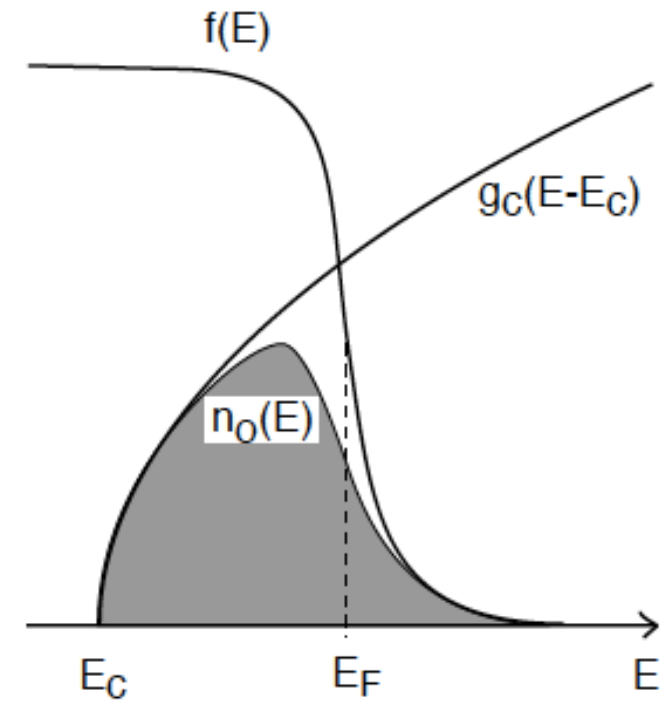
$$n_o p_o = n_i^2$$



Equilibrium **np product** in a semiconductor at a certain temperature is a **constant specific to the semiconductor**.

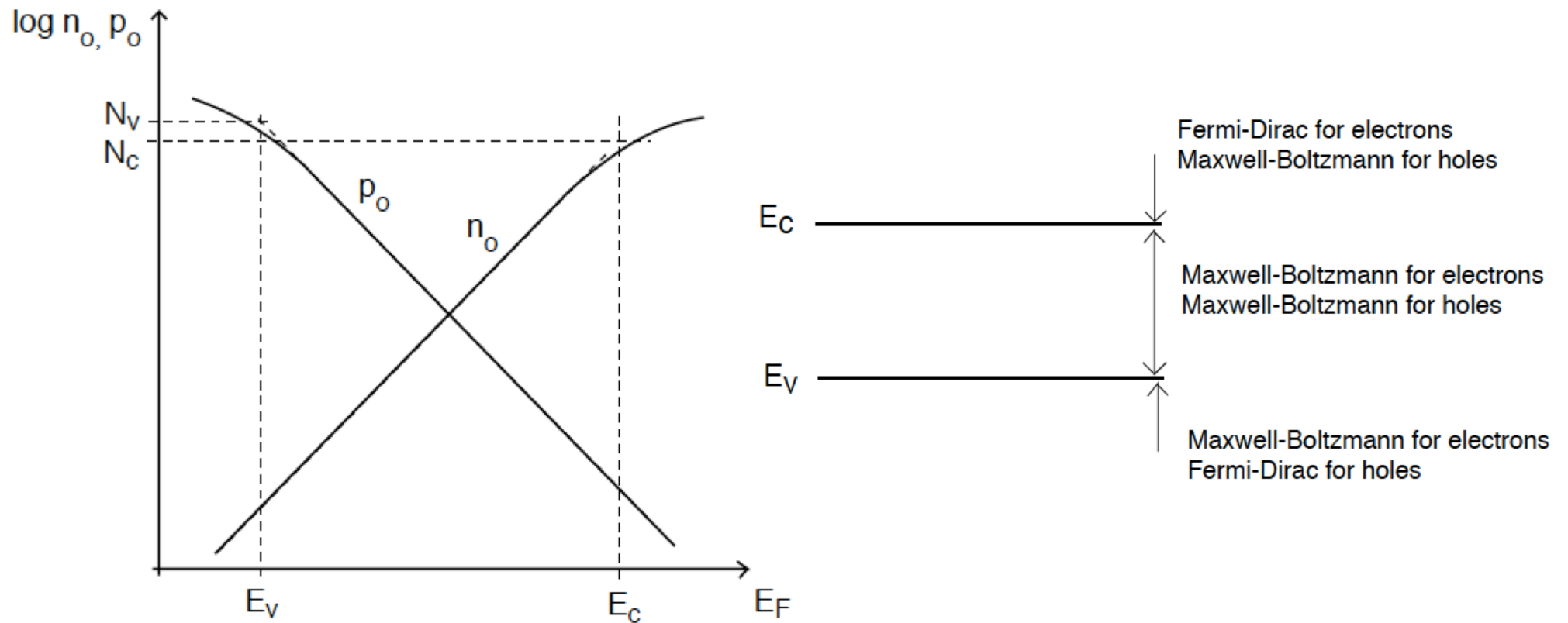


Non degenerate

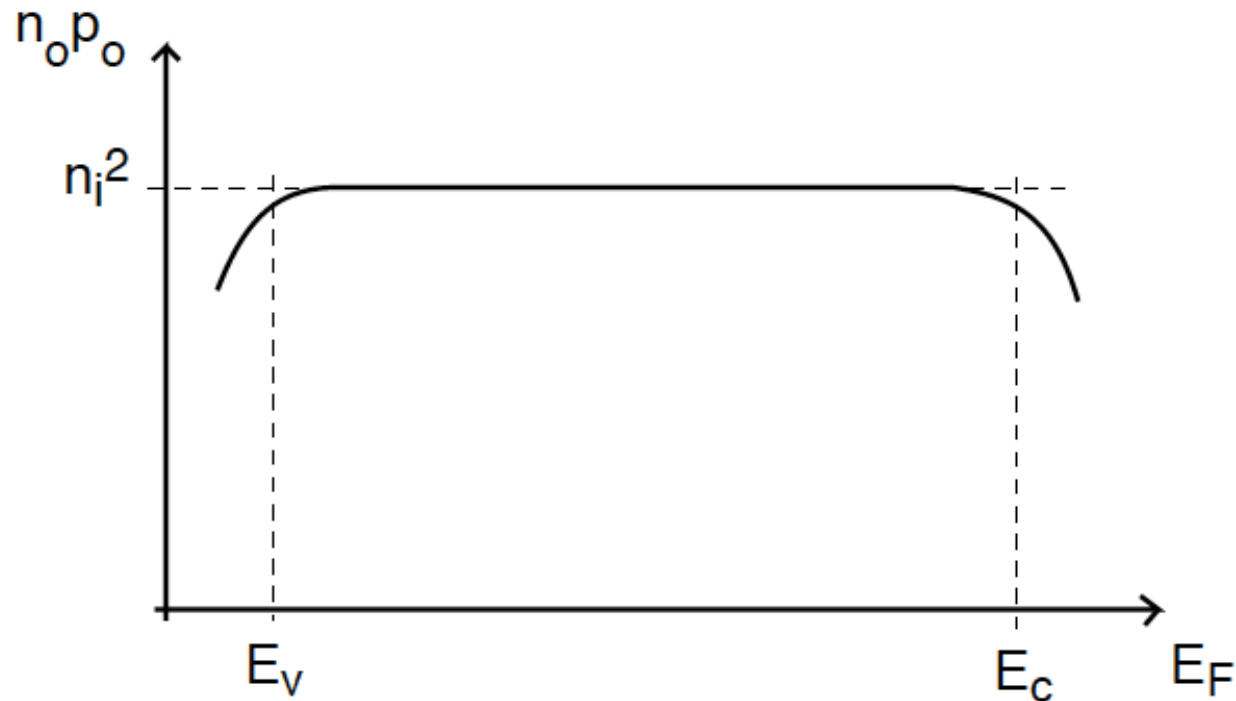


degenerate

Summary of carrier statistics depending on E_F location



Summary of carrier statistics depending on E_F location



If E_F is inside bandgap:

$$n_o p_o \simeq N_c N_v \exp -\frac{E_g}{kT}$$

For a given semiconductor, $n_o p_o$ depends only on T and is independent of precise location of E_F . But only if semiconductor is non-degenerate.

In intrinsic semiconductor, $n_o = p_o$ and usually fairly small $\Rightarrow$ semiconductor non-degenerate.

$$n_i = \sqrt{n_o p_o} = \sqrt{N_c N_v} \exp -\frac{E_g}{2kT}$$

Key conclusions

In solids, electron states cluster in bands separated by bandgaps.

Band structure:

- Depends on the atomic and crystal structure
- Dependent on the momentum of electrons
- Band gaps can be direct or indirect

$E(k)$ is periodic in **k-space**: $E(k) = E(k+G)$

Entire band structures need only to be calculated within the brillouin zone.

Order of magnitude of key parameters:

– Atomic density of Si: $n \sim 5 \times 10^{22} \text{ cm}^{-3}$

– Bandgaps at 300K	Si:	$E_g = 1.12 \text{ eV}$
	GaAs:	$E_g = 1.43 \text{ eV}$
	GaN:	$E_g = 3.39 \text{ eV}$

thermal energy: $kT \sim 26 \text{ meV @300K}$

Distinct feature of semiconductors: at 0 K, quantum state filling ends up with full band separated from next empty band by $\sim 1\text{-}3\text{eV}$ bandgap at around 300 K, some electrons populate next band above bandgap.

System in thermal equilibrium:

- isolated from outside world and in steady state.
- In thermal equilibrium, E_F is independent of position: it is constant!

Occupation probability of quantum systems in thermal equilibrium governed by Fermi-Dirac distribution function:

$$f(E) = \frac{1}{1 + \exp \frac{E - E_F}{kT}}$$

Density of states give us the number of states per volume available at a give energy

In 3D:
$$g_c(E) = 4\pi \left(\frac{2m_{de}^*}{h^2} \right)^{3/2} \sqrt{E - E_c} \quad E \geq E_c$$

$$g_v(E) = 4\pi \left(\frac{2m_{dh}^*}{h^2} \right)^{3/2} \sqrt{E_v - E} \quad E \leq E_v$$

Key conclusions

Non-degenerate semiconductor:

$$n_o = N_c \exp \frac{E_F - E_c}{kT}, \quad p_o = N_v \exp \frac{E_v - E_F}{kT}$$

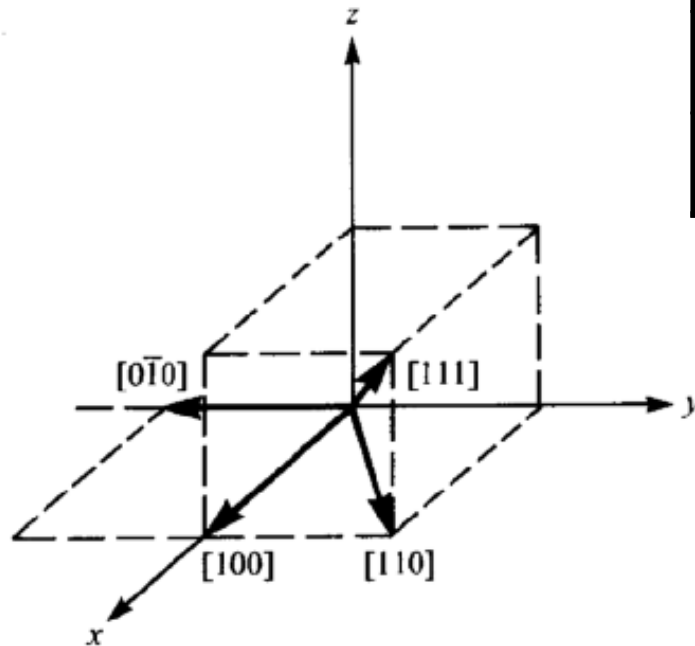
Intrinsic semiconductor: ideally pure semiconductor.

$$n_o = p_o = n_i = \sqrt{N_c N_v} \exp -\frac{E_g}{2kT}$$

In non-degenerate semiconductor $n_o p_o = n_i^2$

Appendix – Supplementary Material

Miller Indices:



Notation	Interpretation
$(h\ k\ l)$	crystal plane
$\{h\ k\ l\}$	equivalent planes
$[h\ k\ l]$	crystal direction
$\langle h\ k\ l \rangle$	equivalent directions

h : inverse x-intercept of plane

k : inverse y-intercept of plane

l : inverse z-intercept of plane

(Intercept values are in multiples of the lattice constant;
 h , k and l are reduced to 3 integers having the same ratio.)

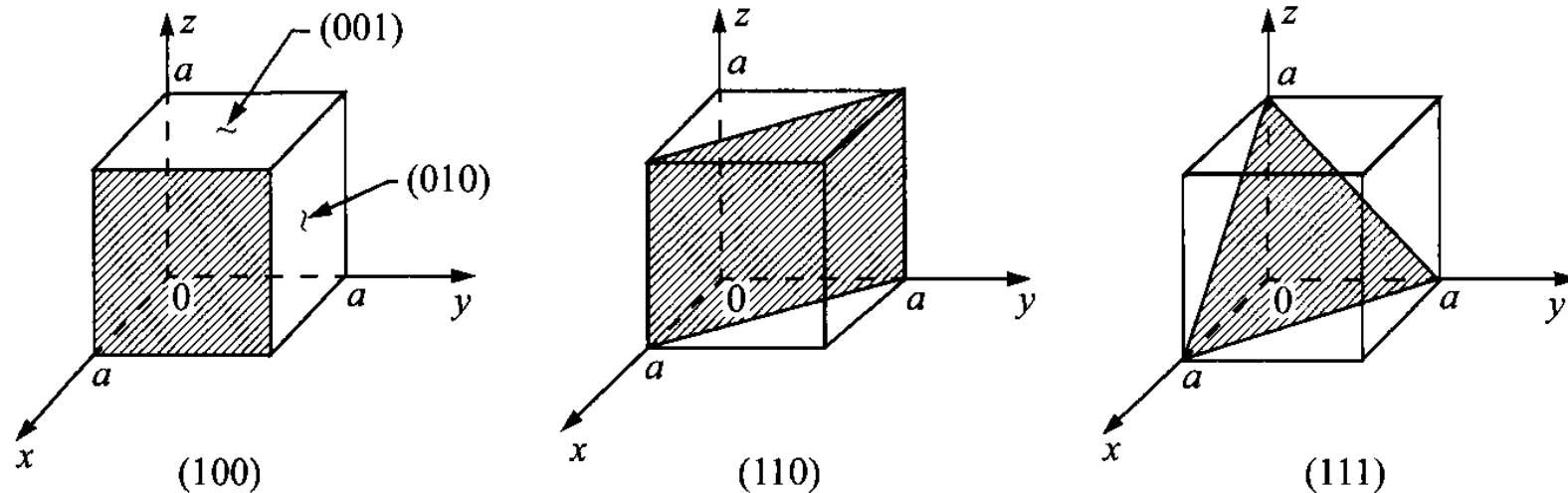
Sample direction vectors and their corresponding Miller indices.

** Wurtzite materials have an hexagonal crystal structure:

4 Miller indices are used to represent it more easily (see appendix for more information)

Crystallographic Notation

Real space or direct space (represented by Miller indices)



Reciprocal space is also called **Fourier space, k- space, or momentum space**

$$\mathbf{a}^* \equiv 2\pi \frac{\mathbf{b} \times \mathbf{c}}{\mathbf{a} \cdot \mathbf{b} \times \mathbf{c}}$$

$$\mathbf{b}^* \equiv 2\pi \frac{\mathbf{c} \times \mathbf{a}}{\mathbf{a} \cdot \mathbf{b} \times \mathbf{c}}$$

$$\mathbf{c}^* \equiv 2\pi \frac{\mathbf{a} \times \mathbf{b}}{\mathbf{a} \cdot \mathbf{b} \times \mathbf{c}}$$

The reciprocal lattice plays a fundamental role in most analytic studies (properties related to momentum) of periodic structures:

- Energy-momentum relationship
- Theory of diffraction: X-Ray diffraction

This is not covered in this course (solid-state physics classes will cover this in details)

Reciprocal lattice vector:

$$\mathbf{G} = h\mathbf{a}^* + k\mathbf{b}^* + l\mathbf{c}^* \quad \text{for which: } \mathbf{G} \cdot \mathbf{R} = 2\pi \times \text{Integer}$$

Crystallographic Notation

Reciprocal space is also called **Fourier space, k- space, or momentum space**

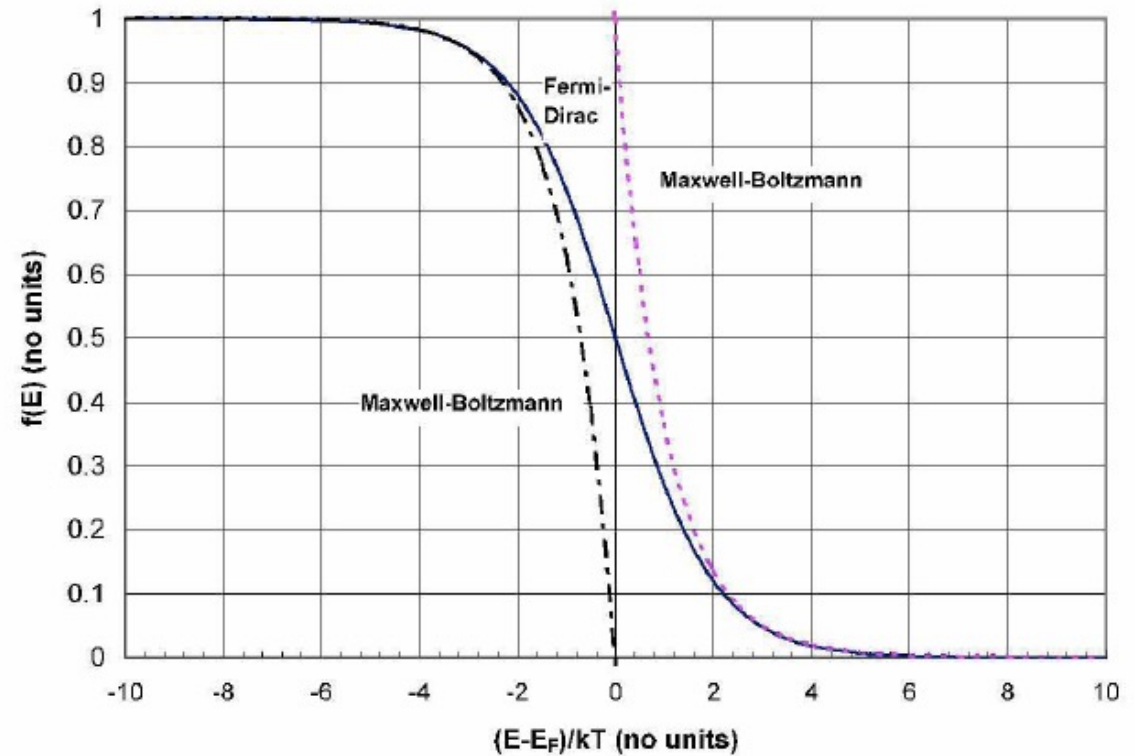
For instance, X-ray diffraction reveals the reciprocal space

Interesting video: <https://www.youtube.com/watch?v=DFFU39A3fPY>

Electron statistics

Approximation

$$f(E) = \frac{1}{1 + \exp \frac{E - E_F}{kT}}$$



- Maxwell-Boltzmann approximation:

For $E - E_F \gg kT$:

$$f(E) \simeq \exp -\frac{E - E_F}{kT}$$

For $E - E_F \ll kT$:

$$f(E) \simeq 1 - \exp \frac{E - E_F}{kT}$$